



# Critical electron scattering in $\text{UGe}_2$ near the magnetic phase transition induced by pressure

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## Abstract

The electrical resistivity of single crystalline ferromagnetic  $\text{UGe}_2$  has been measured under high pressure up to 2 GPa. It is found that the Curie temperature  $T_C$  decreases with increasing pressure and the ferromagnetism disappears at high pressure. The temperature dependent electrical resistivity is analyzed on the basis of spin fluctuation theory. The  $T^2$  term was observed at low temperature and the coefficient shows a divergence near the critical pressure. A  $T^{5/3}$  dependence in the  $\rho(T)$  curve was observed below and above  $T_C$  and its pressure effect is discussed briefly. © 1998 Published by Elsevier Science S.A.

**Keywords:** Electrical resistivity; Ferromagnetism; High pressure; Spin fluctuation

## 1. Introduction

Intermetallic compounds including Ce or U atoms have been investigated extensively because these compounds give important information for studying the role of strong electron correlations in metallic systems [1,2]. It is well known that the physical properties of these compounds having unstable f electrons are strongly affected by external forces such as pressure, magnetic field, etc. [3]. In other words, we can obtain important data for clarifying the electronic structure of these compounds through measurements taken under the influence of these external forces.

$\text{UGe}_2$  crystallizes in the orthorhombic structure with lattice parameters  $a=4.09$  Å,  $b=15.20$  Å and  $c=3.96$  Å [4]. It has been reported that  $\text{UGe}_2$  is a ferromagnet with  $T_C=52$  K and its physical properties show a large anisotropy [4,5]. Recently, a highly qualified single crystal of  $\text{UGe}_2$  was successfully prepared and exhibited an extremely large residual resistivity ratio of several hundreds.

In the present work we made an attempt to measure the temperature dependent electrical resistivity  $\rho(T)$  of single crystalline  $\text{UGe}_2$  at high pressure up to 2 GPa. The  $\rho(T)$  curves are analyzed on the basis of spin fluctuation theory and the critical behavior in  $\rho(T)$  near the phase boundary

is discussed briefly. From the theory of critical phenomena, the similarity between the  $\rho(T)$  curve and the thermal expansion coefficient is pointed out.

## 2. Experimental procedure

Single crystalline  $\text{UGe}_2$  was grown by a Czochralski pulling method. The details of sample preparation and characterization have been reported elsewhere [4]. The electrical resistance along the  $c$ -axis was measured by using a standard d.c. four-probe method with an accuracy of  $10^{-6}$ , in which four gold leads were attached to the sample by means of silver paste. High pressure up to 2 GPa was generated by using a tungsten carbide piston and a Cu–Be cylinder. A 1:1 mixture of Fluorinert FC 70 and FC 77 was used as the pressure transmitting medium. Details of the present high pressure apparatus have been reported previously [6].

## 3. Results and discussion

Fig. 1 shows an example of the electrical resistivity  $\rho$  along the  $c$ -axis ( $J||c$ ) at various pressures as a function of temperature.  $\rho(T)$  at ambient pressure shows a smooth decrease with decreasing temperature and then a sudden decrease is observed near 52 K, which corresponds to the

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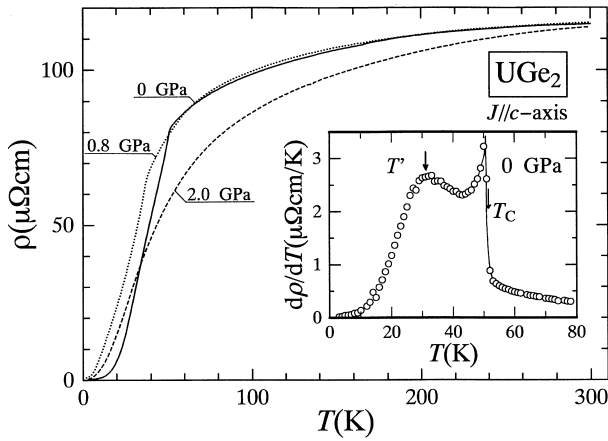


Fig. 1. Temperature dependence of the electrical resistivity  $\rho$  of  $\text{UGe}_2$  along the  $c$ -axis at various pressures. The Curie temperature  $T_C$  and the phase transition temperature  $T'$  at ambient pressure are shown in the inset. Details of the determination of  $T_C$  and  $T'$  are described in the text.

ferromagnetic ordering. The residual resistivity ratio was about 400, which certifies the high quality of the sample. The inset in Fig. 1 shows the value of  $d\rho/dT$  as a function of  $T$  at ambient pressure. There are two peaks in the temperature dependence of  $d\rho/dT$ , a sharp peak due to ferromagnetic ordering and a broad peak which may correspond to a kind of magnetic phase transition, but the origin is not clear at present. The Curie temperature  $T_C$  is defined as the temperature of the half value of  $(d\rho/dT)_{\text{peak}} - (d\rho/dT)_{\text{para}}$ , which is shown by an arrow in the inset. The transition temperature  $T'$  corresponding to the broad peak is determined as the temperature where  $d\rho/dT$  begins to decrease with decreasing temperature, which is also shown in the inset. An anomaly near  $T'$  has been observed in the thermal expansion coefficient, the Hall coefficient and the thermoelectric power [4,5,7]. It is suggested that the ferromagnetism of  $\text{UGe}_2$  is not simple because the  $\rho(T)$  curve along the  $b$ -axis shows a hump below  $T_C$  which reminds us of antiferromagnetism with a gap [4].  $T_C$  and  $T'$  are shown in Fig. 2 as a function of pressure. Both  $T_C$  and  $T'$  were found to decrease by

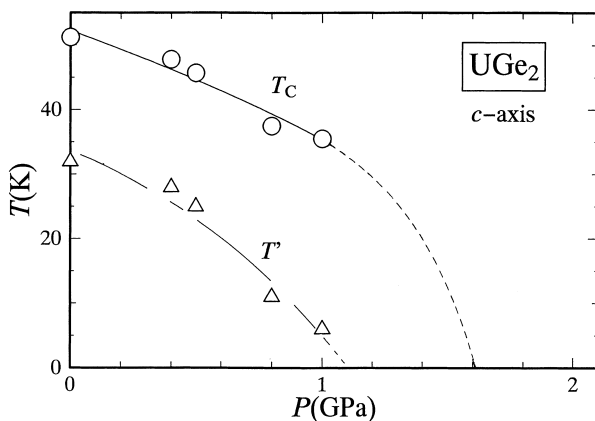


Fig. 2. The critical temperatures  $T_C$  and  $T'$  as a function of pressure.

applying pressure with coefficients  $\partial T_C/\partial P = -14 \text{ K GPa}^{-1}$  and  $\partial T'/\partial P = -23 \text{ K GPa}^{-1}$ , respectively. The anomaly at  $T_C$  in the  $\rho(T)$  curve due to ferromagnetic ordering disappeared above 1.5 GPa. The solid (partly dotted) curve for  $T_C$  and  $T'$  are to guide the eye. Above 1.5 GPa,  $\text{UGe}_2$  is considered to be nearly ferromagnetic or a highly correlated compound, in which the scattering of conduction electrons is dominated mainly by spin fluctuation [4].

According to the self-consistent renormalized spin fluctuation theory of Ueda and Moriya (UM) [8], the electrical resistivity  $\rho(T)$  of a highly correlated itinerant ferromagnet shows a temperature dependence of  $T^2$  in the temperature range  $T \ll T_C$  and of  $T^{5/3}$  both in the ranges  $T \geq T_C$  and  $T \leq T_C$ . Since the 5f electrons in  $\text{UGe}_2$  are considered to be strongly correlated itinerant electrons with spin fluctuations like 3d transition metals [9], the UM theory may be applicable to explain the present results.

In order to examine the  $T^2$  dependence in the  $\rho(T)$  curve due to spin fluctuation and spin wave excitation in the range  $T \ll T_C$  we plotted  $\rho$  at low temperature as a function of  $T^2$ , which is shown in Fig. 3. A  $T^2$  dependence is observed in this plot. The range showing a  $T^2$  dependence depends on the pressure. The coefficients  $A(P)$  in the equation  $\rho = \rho_0 + AT^2$ , where  $\rho_0$  is the residual resistivity, are obtained by calculating the slope of the  $\rho$  vs.  $T^2$  plot in Fig. 3. It is easily seen that the slope of the plot is large near the critical boundary of ferromagnetism. The values of  $A$  are shown in Fig. 4 as a function of pressure.  $A$  is found to increase with pressure below 1 GPa followed by a decrease above 1.5 GPa. In other words,  $A$  shows a divergence at a critical pressure  $P_c$ . The mag-

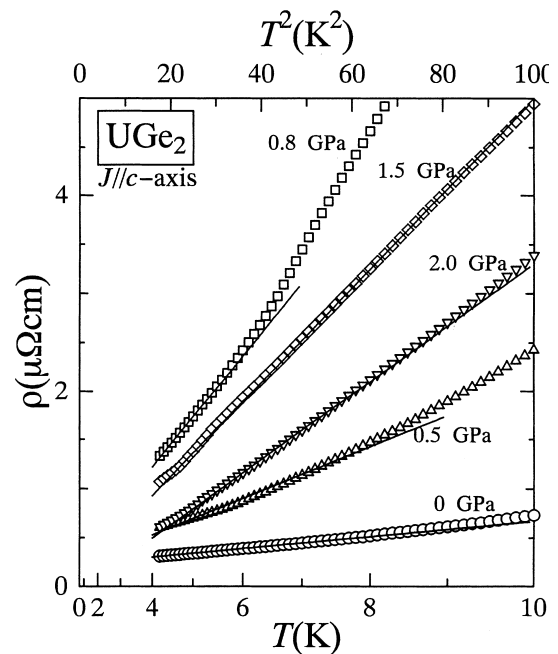


Fig. 3.  $\rho$  at low temperature in quadratic scale. The solid lines show the  $T^2$  dependence.

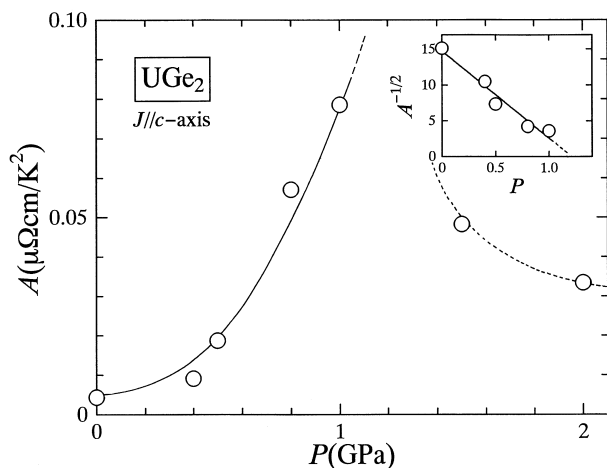


Fig. 4. The coefficient of the  $T^2$  term as a function of pressure. The solid and dashed curves are a guide to the eye. The pressure dependence of the value of  $1/A^{-1/2}$ , which corresponds to the characteristic energy, is plotted in the inset.

nitude of  $A$  is known to be proportional to  $T_0^{-2}$  [3], where  $T_0$  is a characteristic spin fluctuation temperature, i.e.  $T_0 \propto A^{-1/2}$ . The pressure dependence of  $A^{-1/2}$  is shown in the inset of Fig. 4.  $A^{-1/2}$  or  $T_0$  is extrapolated to 0 around 1.3 GPa. In the following we assume  $P_c \approx 1.3$  GPa.

It has been reported that the value of  $A$  shows a divergence near the critical point (concentration or pressure as control parameter) where the magnetic ordering disappears or begins [10,11]. It was reported by Takahashi et al. [12] that the resistance anomaly due to ferromagnetic ordering disappears around 1.6 GPa. Considering these facts, the divergence observed in the pressure dependence of  $A$  may correspond to the disappearance of the ferromagnetism of  $\text{UGe}_2$  around 1.3 GPa and the effect of spin fluctuation on the  $\rho(T)$  curve is the most significant.

Here we examine the temperature dependence of  $\rho$  for  $\text{UGe}_2$  at high pressure. Fig. 5 shows  $\rho$  as a function of reduced temperature  $(T/T_c)^{5/3}$  at various pressures up to 1 GPa in the temperature range  $0.5 \leq (T/T_c)^{5/3} \leq 1.6$ . At ambient pressure the  $T^{5/3}$  temperature dependence is

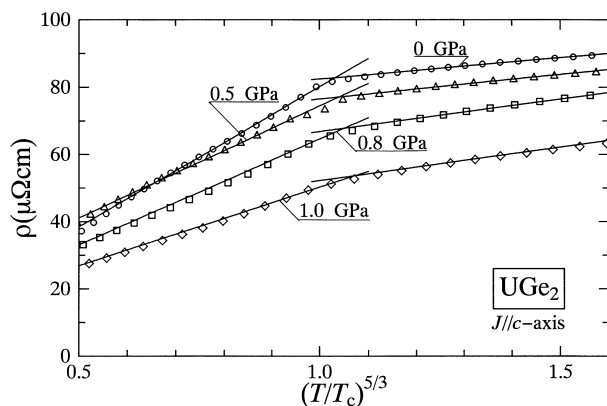


Fig. 5.  $\rho$  plotted as a function of  $T^{5/3}$  around  $T_c$ . The solid lines show the  $T^{5/3}$  dependence both above and below  $T_c$ .

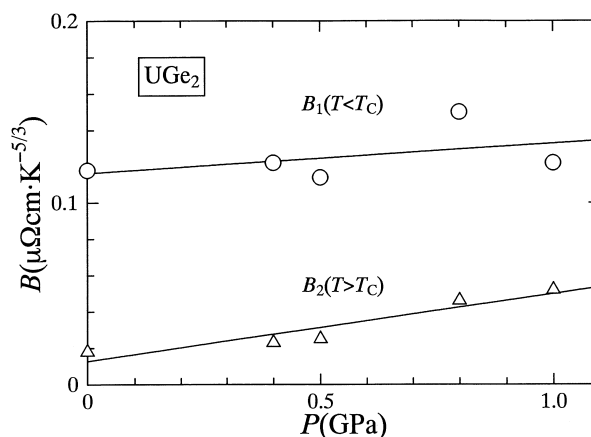


Fig. 6. Pressure dependence of the coefficient of the  $T^{5/3}$  term  $B_1$  obtained below  $T_c$  and  $B_2$  above  $T_c$ .

clearly seen, as shown by the solid line used to guide the eye. As the pressure increases, the range showing a  $T^{5/3}$  dependence becomes broader. The coefficients of  $T^{5/3}$ ,  $B_1$  below  $T_c$  and  $B_2$  above  $T_c$ , are shown in Fig. 6. The values of  $B_1$  and  $B_2$  roughly increase with pressure: the increase in the magnitude of  $B_2$  is larger than that of  $B_1$ . In the limit of critical fluctuations,  $B_1$  is expected to be equal to  $B_2$  [8], i.e.  $B_1/B_2 = 1$ . In Fig. 7 we plot  $B_1/B_2$  as a function of pressure. It is found that  $B_1/B_2$  extrapolates to 1 around 1.3 GPa, which is in agreement with the critical pressure  $P_c$  obtained above. This supports that the above consideration on the basis of spin fluctuation theory is reasonable even for  $\text{UGe}_2$ . But the value of  $P_c$  ( $\approx 1.3$  GPa) in the present work is different from that obtained with thermal expansion measurements ( $P_c \approx 1.8$  GPa) [7]. As was mentioned previously, there may be another kind of phase transition below  $T'$  ( $\approx 30$  K at ambient pressure) which has been reported to decrease with pressure and disappear near 1 GPa [13]. For example, at 5 K, using the results of Fig. 2,  $\text{UGe}_2$  is ferromagnetic at ambient

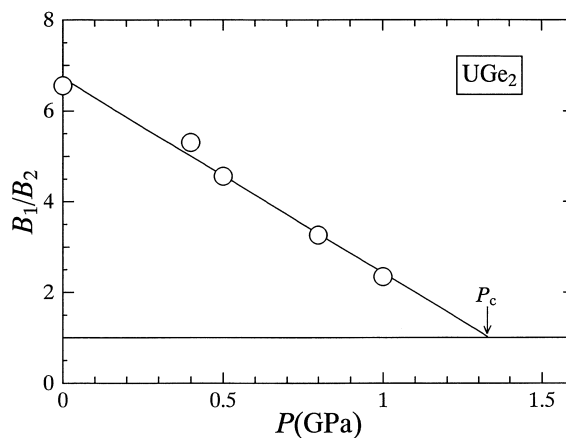


Fig. 7. The ratio of  $B_1$  and  $B_2$  as a function of pressure. The result of a least-squares fit is shown as the solid line and the critical pressure  $P_c$  is determined as the pressure at which it extrapolates to 1.

pressure, but as the pressure increases it undergoes a phase transition around 1 GPa, which may be a change in spin structure of second order, and then finally it becomes paramagnetic or nearly ferromagnetic above 1.8 GPa. For this reason, the phase transition at  $T'$  may make the determination of  $P_c$  difficult by electrical resistance measurements. In other words, the scattering mechanism of the conduction electrons below  $T'$  is considered to be very complicated, as suggested in the following paragraph. Thus the electronic and magnetic structure of  $\text{UGe}_2$  in the pressure range between 1 and 1.8 GPa seems not to be simple, as pointed out previously [9].

Finally, we should mention the similarity between the critical behavior in the electrical resistivity and thermal expansion coefficients of  $\text{UGe}_2$ . It is well known that the magnetic contribution to the temperature derivative of the electrical resistivity  $d\rho/dT$  should be proportional to the magnetic part of the specific heat  $C$ ,  $d\rho/dT \propto C$  [14]. Furthermore, the Grüneisen relation is described as

$$\alpha(T) = \frac{\kappa\Gamma}{3V}C \quad (1)$$

where  $\kappa$ ,  $V$  and  $\Gamma$  are the compressibility, the volume and the Grüneisen parameter, respectively. Since  $\kappa$ ,  $V$  and  $\Gamma$  are weakly dependent on temperature, we can consider that  $\alpha$  is roughly proportional to  $C$ . Then  $\alpha$  is also proportional to  $d\rho/dT$ . Here we apply this consideration to the present results. Fig. 8 shows the temperature dependence of  $\alpha_c$

$[=(1/c)(dc/dT)]$  and  $d\rho/dT$ . We can easily find the similarity between these quantities. By applying pressure,  $T_c$  decreases and the width of the peak due to ferromagnetic ordering becomes broad. It should be noted that the shape of  $d\rho/dT$  at 1 GPa is significantly different from that at ambient pressure, but in the case of  $\alpha_c(T)$  there is no such large difference. This suggests that the conduction electrons are scattered in a complicated way below  $T_c$ . The paramagnetic state is induced at 2 GPa and there are no clear peaks in the temperature dependence of  $\alpha_c$  and  $\rho$  except a broad peak centered near 27 K in the  $\alpha_c(T)$  curve. Such a similarity has been reported for other rare earth compounds [15].

#### 4. Conclusion

We have observed the electrical resistance of single crystalline  $\text{UGe}_2$  along the  $c$ -axis. The ferromagnetism disappeared at high pressure around  $P_c = 1.3$  GPa. At  $P_c$ , the coefficient of the  $T^2$  term is found to show a divergence. The critical behavior near  $T_c$  is explained using the self-consistent renormalized spin fluctuation theory. The similarity between the thermal expansion coefficients and  $d\rho/dT$  exists not only at ambient pressure but also at high pressure.

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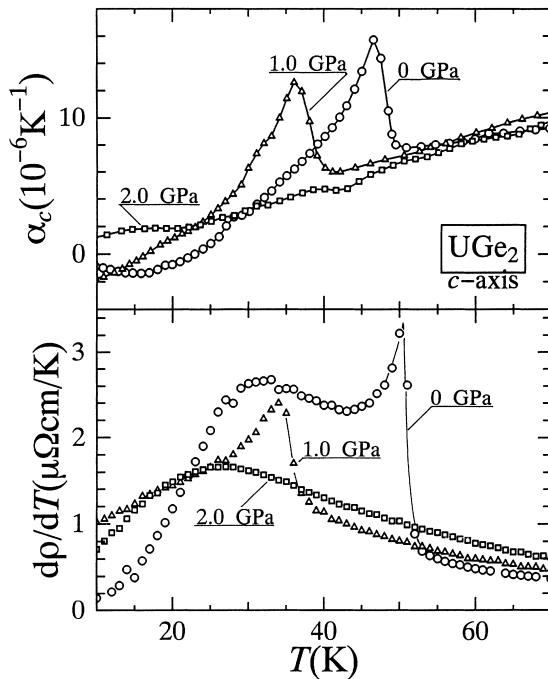


Fig. 8. Temperature dependence of (a) the thermal expansion coefficient  $\alpha_c$  and (b) the temperature derivative of  $\rho$  at high pressure.